# Perfect state transfer on Cayley graphs over a group of order $8 n$ 

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#### Abstract

The transition matrix of a graph $\Gamma$ with adjacency matrix $A$ is defined by $H(t):=\exp (-\mathbf{i} t A)$, where $t \in \mathbb{R}$ and $\mathbf{i}=\sqrt{-1}$. The graph $\Gamma$ is said to admit perfect state transfer (PST) between a pair of vertices $u$ and $v$ if there exists $\tau(>0) \in \mathbb{R}$ such that $\left|H(\tau)_{u v}\right|=1$. Perfect state transfer has great importance due to its applications in quantum information processing, quantum communication networks and cryptography. In this paper, we study the existence of perfect state transfer on Cayley graphs over the group $V_{8 n}$. we present some necessary and sufficient conditions for the existence of perfect state transfer on Cay $\left(V_{8 n}, S\right)$.


Keywords. Perfect state transfer; Cayley graph; Eigenvalues of a graph
Mathematics Subject Classifications: 05C25; 81P45; 81Q35

## 1 Introduction

The concept of quantum walks on graphs was first introduced by Farhi and Gutmann [14] in the year 1998. Quantum walks on finite graphs provide useful simple models for quantum transport phenomena which was first discovered by Bose [7] in 2003. Christandl et al. [10] proposed a class of qubit networks that admit perfect state transfer. Quantum walks are important tools in quantum computation and information theory and can be used to describe the fidelity of information transfer in a network of interacting qubits.

Let $\Gamma$ be a finite simple connected graph with adjacency matrix $A$. Denoted by $V(\Gamma)$ the set of vertices of $\Gamma$. The transition matrix of $\Gamma$ is defined by

$$
H(t)=H_{\Gamma}(t)=: \exp (-\mathbf{i} t A)=\sum_{s=0}^{\infty} \frac{(-\mathbf{i} t A)^{s}}{s!}=\left(H_{u, v}(t)\right)_{u, v \in V(\Gamma)},
$$

where $t \in \mathbb{R}$ and $\mathbf{i}=\sqrt{-1}$.

A graph $\Gamma$ is said to exhibit perfect state transfer (PST) from a vertex $u$ to the vertex $v$ if there exists $\tau(>0) \in \mathbb{R}$ such that the $u v$-th element of $H(\tau)$ has absolute value 1 . We describe $\Gamma$ as exhibiting periodicity at the vertex $u$ at time $\tau$ if the $u u$-th element of $H(\tau)$ has absolute value 1 . The graph $\Gamma$ is considered to exhibit periodicity if it exhibits the property of periodicity across all its vertices simultaneously at the same time.

PST has been studied for several families of graphs. Angeles-Canul et al. [1] investigated PST in integral circulant graphs and the join of graphs. Coutinho and Godsil [11] explored PST in products and covers of graphs. Pal and Bhattacharjya [22] studied PST on NEPS of the path on three vertices. Godsil [17] offers a survey on perfect state transfer and related questions up to 2011. He [16] explains the close relationship between the existence of perfect state transfer on certain graphs and association schemes. Notably, in [18] Godsil presents a complete characterization of PST on simple connected graphs.

Cayley graphs are good candidates for exhibiting PST due to their nice algebraic structure. Among these results Basic et al. [4, 5, 3, 23], Cheung and Godsil [9] and Bernasconi et al. [6] present some criterions on circulant graphs and cubelike graphs having PST. Tan et al. [25] presented a characterization on abelian Cayley graphs having PST. They showed that many of the previous results on periodicity and existence of PST in circulant graphs and cubelike graphs can be derived in unified and more simple ways. However, relatively little research has been carried out on Cayley graphs over non-abelian groups exhibiting PST. Cao and Feng [8] investigated PST on Cayley graphs over dihedral groups. Subsequently, Arezoomand et al. [2] and Luo et al. [21] explored PST on Cayley graphs over dicyclic groups and semidihedral groups, respectively. Recently, Khalilipour and Ghorbani [20] studied PST on Cayley graphs over the group $U_{6 n}=\left\langle a, b: a^{2 n}=b^{3}=1, a^{-1} b a=b^{-1}\right\rangle$.

In this paper, we consider the existence of PST on Cayley graphs over the group $V_{8 n}=\left\langle a, b: a^{2 n}=\right.$ $\left.b^{4}=1, \quad b a=a^{-1} b^{-1}, b^{-1} a=a^{-1} b\right\rangle$, where $n$ is a positive integer. Using the irreducible representations of $V_{8 n}$, several necessary and sufficient conditions for a normal Cayley graph Cay $\left(V_{8 n}, S\right)$ exhibiting PST are carried out.

The rest of the current work is organized as follows. In Section 2, we give the description of the irreducible representations of $V_{8 n}$ and spectra of normal Cayley graphs Cay $\left(V_{8 n}, S\right)$. The existence of PST on normal Cayley graphs $\operatorname{Cay}\left(V_{8 n}, S\right)$ is explored in Section 3.

## 2 Irreducible representations of the group $V_{8 n}$ and spectra of Cayley graphs $\operatorname{Cay}\left(V_{8 n}, S\right)$

A representation of a finite group $G$ is a homomorphism $\theta: G \rightarrow G L(U)$, where $G L(U)$ is the group of all automorphisms of a finite-dimensional and non-zero complex vector space $U$. The dimension of $U$
is called the degree of $\theta$. Two representations $\theta$ and $\psi$ of $G$ on $U$ and $W$, respectively, are equivalent, denoted by $\theta \sim \psi$, if there is an isomorphism $T: U \rightarrow W$ such that $\theta(g)=T \psi(g) T^{-1}$ for all $g \in G$.

Let $\theta: G \rightarrow G L(U)$ be a representation. The character $\chi_{\theta}: G \rightarrow \mathbb{C}$ of $\theta$ is defined by setting $\chi_{\theta}(g)=\operatorname{Tr}(\theta(g))$ for all $g \in G$, where $\operatorname{Tr}(\theta(g))$ is the trace of the representation matrix of $\theta(g)$. A subspace $W$ of $U$ is said to be $G$-invariant if $\theta(g) w \in W$ for all $g \in G$ and $w \in W$. Obviously, $\{0\}$ and $U$ are $G$-invariant subspaces, called trivial subspaces. If $U$ has no non-trivial $G$-invariant subspaces, then $\theta$ is called an irreducible representation and $\chi_{\theta}$ an irreducible character of $G$.

Let $G$ be a finite group and $S$ be a symmetric subset of $G$, that is, $S=S^{-1}$, where $S^{-1}=\left\{s^{-1}: s \in S\right\}$ and $\mathbf{1} \notin S$. The Cayley graph of $G$ with respect to $S$, denoted Cay $(G, S)$, is a graph whose vertices are the elements of $G$ and there exists an edge between distinct vertices $g, h \in G$ if $g h^{-1} \in S$. The set $S$ is called the connection set. If $S g=g S$ for all $g \in G$, then $S$ is called a normal Cayley subset and Cay $(G, S)$ a normal Cayley graph. Since $S$ is symmetric, $\operatorname{Cay}(G, S)$ is a simple graph. We assume $G=\langle S\rangle$ to ensure that $\operatorname{Cay}(G, S)$ is a connected graph. The adjacency matrix of $\operatorname{Cay}(G, S)$ is defined by $A=\left(a_{g, h}\right)_{g, h \in G}$, where

$$
a_{g, h}= \begin{cases}1 & \text { if } g h^{-1} \in S \\ 0 & \text { otherwise }\end{cases}
$$

For more properties about Cayley graphs, One can refer to [15].
Let $n$ be a positive integer. The group $V_{8 n}$ is defined by

$$
V_{8 n}=\left\langle a, b: a^{2 n}=b^{4}=1, \quad b a=a^{-1} b^{-1}, b^{-1} a=a^{-1} b\right\rangle .
$$

Note that $V_{8 n}=\left\{a^{r}, a^{r} b, a^{r} b^{2}, a^{r} b^{3}: 0 \leq r \leq 2 n-1\right\}$
For odd values of $n$, the group $V_{8 n}$ has $2 n+3$ conjugacy classes as follows:

$$
\begin{array}{ll}
\{1\},\left\{b^{2}\right\},\left\{a^{2 r+1}, a^{-2 r-1} b^{2}\right\}, & r \in\{0, \ldots, n-1\}, \\
\left\{a^{2 s}, a^{-2 s}\right\},\left\{a^{2 s} b^{2}, a^{-2 s} b^{2}\right\}, & s \in\{1, \ldots,(n-1) / 2\}, \\
\left\{a^{j} b^{k}: j \text { even, } k=1 \text { or } 3\right\} \text { and } \\
\left\{a^{j} b^{k}: j \text { odd, } k=1 \text { or } 3\right\} .
\end{array}
$$

For even values of $n$, the group $V_{8 n}$ has $2 n+6$ conjugacy classes as follows:

$$
\begin{aligned}
& \{1\},\left\{b^{2}\right\},\left\{a^{n}\right\},\left\{a^{n} b^{2}\right\}, \\
& \left\{a^{2 r+1}, a^{-(2 r+1)} b^{2}\right\}, \quad r \in\{0, \ldots, n-1\} \\
& \left\{a^{2 s}, a^{-2 s}\right\},\left\{a^{2 s} b^{2}, a^{-2 s} b^{2}\right\}, \quad s \in\{1, \ldots, n / 2-1\} \\
& \left\{a^{2 k} b^{(-1)^{k}}: 0 \leq k \leq n-1\right\}, \quad\left\{a^{2 k} b^{(-1)^{k+1}}: 0 \leq k \leq n-1\right\} \\
& \left\{a^{2 k+1} b^{(-1)^{k}}: 0 \leq k \leq n-1\right\} \text { and }\left\{a^{2 k+1} b^{(-1)^{k+1}}: 0 \leq k \leq n-1\right\}
\end{aligned}
$$

Table 1: Irreducible representations of $V_{8 n}$, for $n$ odd

|  | $a$ | $b$ |
| :---: | :---: | :---: |
| $\chi_{1}$ | 1 | 1 |
| $\chi_{2}$ | 1 | -1 |
| $\chi_{3}$ | -1 | 1 |
| $\chi_{4}$ | -1 | -1 |
| $\psi_{j}(0 \leq j \leq n-1)$ | $\left(\begin{array}{cc}\omega^{2 j} & 0 \\ 0 & -\omega^{-2 j}\end{array}\right)$ | $\left(\begin{array}{cc}0 & 1 \\ -1 & 0\end{array}\right)$ |
| $\phi_{k}(1 \leq k \leq n-1)$ | $\left(\begin{array}{cc}\omega^{k} & 0 \\ 0 & \omega^{-k}\end{array}\right)$ | $\left(\begin{array}{ll}0 & 1 \\ 1 & 0\end{array}\right)$ |

The following lemma presents the well-known irreducible representations and characters for the group $V_{8 n}$.

Lemma 2.1. [19, 13] Let $n$ be a positive integer and $\omega=\exp \left(\frac{2 \pi i}{2 n}\right)$ be a primitive $2 n$-th root of unity.
(1) The irreducible representations of $V_{8 n}$ are listed in Table 1 for $n$ odd and in Table 2 for $n$ even.
(2) The character table of $V_{8 n}$ is listed in Table 3 for $n$ odd, in Table 4 for $n \equiv 0(\bmod 4)$ and in Table 5 for $n \equiv 2(\bmod 4)$.

The following lemma determines the eigenvalues and eigenvectors of the adjacency matrix of a normal Cayley graph.

Lemma 2.2. [24] Let $G=\left\{g_{1}, \ldots, g_{n}\right\}$ be a finite group and $\phi^{(1)}, \ldots, \phi^{(t)}$ be a complete set of unitary representatives of the equivalence classes of irreducible representations of $G$. Let $\chi_{k}$ be the character of $\phi^{(k)}$ and $d_{k}$ be the degree of $\phi^{(k)}$. Let $S$ be a symmetric subset of $G$ and assume further that $S$ is conjugation-closed. Then the eigenvalues of the Cayley graph Cay $(G, S)$ are $\lambda_{1}, \ldots, \lambda_{t}$, where

$$
\lambda_{k}=\frac{1}{d_{k}} \sum_{s \in S} \chi_{k}(s), \quad 1 \leq k \leq t
$$

and $\lambda_{k}$ has multiplicity $d_{k}^{2}$. Moreover, the vectors

$$
v_{i j}^{(k)}=\sqrt{\frac{d_{k}}{|G|}}\left(\phi_{i j}^{(k)}\left(g_{1}\right), \ldots, \phi_{i j}^{(k)}\left(g_{n}\right)\right)^{t}, 1 \leq i, j \leq d_{k}
$$

form an orthonormal basis for the eigenspace associated with the eigenvalue $\lambda_{k}$.

Table 2: Irreducible representations of $V_{8 n}$, for $n$ even

|  | $a$ | $b$ |
| :---: | :---: | :---: |
| $\chi_{1}$ | 1 | 1 |
| $\chi_{2}$ | $\mathbf{i}$ | $-\mathbf{i}$ |
| $\chi_{3}$ | -1 | -1 |
| $\chi_{4}$ | $-\mathbf{i}$ | $\mathbf{i}$ |
| $\chi_{5}$ | 1 | -1 |
| $\chi_{6}$ | $\mathbf{i}$ | $\mathbf{i}$ |
| $\chi_{7}$ | -1 | 1 |
| $\chi_{8}$ | $-\mathbf{i}$ | $-\mathbf{i}$ |
| $\psi_{j}(1 \leq j \leq n-1)$ | $\left(\begin{array}{cc}\omega^{j} & 0 \\ 0 & \omega^{-j}\end{array}\right)$ | $\left(\begin{array}{cc}0 & \mathbf{i} \\ -\mathbf{i} & 0\end{array}\right)$ |
| $\phi_{k}(1 \leq k \leq n-1)$ | $\left(\begin{array}{cc}\mathbf{i} \omega^{k} & 0 \\ 0 & \mathbf{i} \omega^{-k}\end{array}\right)$ | $\left(\begin{array}{cc}0 & 1 \\ -1 & 0\end{array}\right)$ |

Table 3: Character table of $V_{8 n}$, for $n$ odd

|  | 1 | $b^{2}$ | $a^{2 r+1}(0 \leq r \leq n-1)$ | $a^{2 s}$ | $a^{2 s} b^{2}(1 \leq s \leq(n-1) / 2)$ | $b$ | $a b$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\xi_{1}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\xi_{2}$ | 1 | 1 | 1 | 1 | 1 |  |  |
| $\xi_{3}$ | 1 | 1 | -1 | 1 | 1 | -1 |  |
| $\xi_{4}$ | 1 | 1 | -1 | 1 | 1 | 1 | -1 |
| $\zeta_{j}(0 \leq j \leq n-1)$ | 2 | -2 | $\omega^{2 j(2 r+1)}-\omega^{-2 j(2 r+1)}$ | $\omega^{4 j s}+\omega^{-4 j s}$ | $-\omega^{4 j s}-\omega^{-4 j s}$ | 0 | 0 |
| $\nu_{k}(1 \leq k \leq n-1)$ | 2 | 2 | $\omega^{k(2 r+1)}+\omega^{-k(2 r+1)}$ | $\omega^{2 k s}+\omega^{-2 k s}$ | $\omega^{2 k s}+\omega^{-2 k s}$ | 0 | 0 |

Table 4: Character table of $V_{8 n}$, for $n \equiv 0(\bmod 4)$

|  | 1 | $b^{2}$ | $a^{n}$ | $a^{n} b^{2}$ | $a^{4 m+1}$ | $a^{4 m+3}$ | $a^{4 s}$ | $a^{4 t+2}$ | $a^{4 s} b^{2}$ | $a^{4 t+2} b^{2}$ | $b$ | $b^{-1}$ | $a b$ | $a b^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\xi_{1}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\xi_{2}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 |
| $\xi_{3}$ | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 |
| $\xi_{4}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 |
| $\xi_{5}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 |
| $\xi_{6}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{- 1}$ | -1 | 1 |
| $\xi_{7}$ | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | 1 | 1 | 1 | -1 | -1 |
| $\xi_{8}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | -1 | 1 |
| $\zeta_{j}$ | 2 | 2 | $2(-1)^{j}$ | $2(-1)^{j}$ | $\alpha^{j(4 m+1)}$ | $\alpha^{j(4 m+3)}$ | $\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | $\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | 0 | 0 | 0 | 0 |
| $\nu_{k}$ | 2 | -2 | $2(-1)^{k}$ | $-2(-1)^{k}$ | $\mathbf{i} \alpha^{j(4 m+1)}$ | $\mathbf{- i} \alpha^{j(4 m+3)}$ | $\alpha^{j(4 s)}$ | $-\alpha^{j(4 t+2)}$ | $-\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | 0 | 0 | 0 | 0 |

$\alpha^{j r}=\omega^{j r}+\omega^{-j r}=2 \cos \left(\frac{\pi j r}{n}\right), \alpha^{k r}=\omega^{k r}+\omega^{-k r}=2 \cos \left(\frac{\pi k r}{n}\right), \omega=\exp \left(\frac{2 \pi \mathbf{i}}{2 n}\right) ;$
$m \in\{0, \ldots, n / 2-1\}, s \in\{1, \ldots, n / 4-1\}, t \in\{0, \ldots, n / 4-1\}, j, k \in\{1, \ldots, n-1\}$.

Table 5: Character table of $V_{8 n}$, for $n \equiv 2(\bmod 4)$

|  | 1 | $b^{2}$ | $a^{n}$ | $a^{n} b^{2}$ | $a^{4 m+1}$ | $a^{4 m+3}$ | $a^{4 s}$ | $a^{4 t+2}$ | $a^{4 s} b^{2}$ | $a^{4 t+2} b^{2}$ | $b$ | $b^{-1}$ | $a b$ | $a b^{-1}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\xi_{1}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 |
| $\xi_{2}$ | 1 | -1 | 1 | -1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 |
| $\xi_{3}$ | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 |
| $\xi_{4}$ | 1 | -1 | 1 | -1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 |
| $\xi_{5}$ | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | -1 | -1 | -1 | -1 |
| $\xi_{6}$ | 1 | -1 | 1 | -1 | $\mathbf{i}$ | $-\mathbf{i}$ | 1 | -1 | -1 | 1 | $\mathbf{i}$ | $-\mathbf{i}$ | -1 | 1 |
| $\xi_{7}$ | 1 | 1 | 1 | 1 | -1 | -1 | 1 | 1 | 1 | 1 | 1 | 1 | -1 | -1 |
| $\xi_{8}$ | 1 | -1 | 1 | -1 | $-\mathbf{i}$ | $\mathbf{i}$ | 1 | -1 | -1 | 1 | $-\mathbf{i}$ | $\mathbf{i}$ | -1 | 1 |
| $\zeta_{j}$ | 2 | 2 | $2(-1)^{j}$ | $2(-1)^{j}$ | $\alpha^{j(4 m+1)}$ | $\alpha^{j(4 m+3)}$ | $\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | $\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | 0 | 0 | 0 | 0 |
| $\nu_{k}$ | 2 | -2 | $-2(-1)^{k}$ | $2(-1)^{k}$ | $\mathbf{i} \alpha^{j(4 m+1)}$ | $-\mathbf{i} \alpha^{j(4 m+3)}$ | $\alpha^{j(4 s)}$ | $-\alpha^{j(4 t+2)}$ | $-\alpha^{j(4 s)}$ | $\alpha^{j(4 t+2)}$ | 0 | 0 | 0 | 0 |

$\alpha^{j r}=\omega^{j r}+\omega^{-j r}=2 \cos \left(\frac{\pi j r}{n}\right), \alpha^{k r}=\omega^{k r}+\omega^{-k r}=2 \cos \left(\frac{\pi k r}{n}\right), \omega=\exp \left(\frac{2 \pi \mathbf{i}}{2 n}\right) ;$
$m \in\{0, \ldots, n / 2-1\}, s \in\{1, \ldots, n / 4-1\}, t \in\{0, \ldots, n / 4-1\}, j, k \in\{1, \ldots, n-1\}$.

### 2.1 Spectra of normal Cayley graphs $\operatorname{Cay}\left(V_{8 n}, S\right)$, for odd $n$

We consider the fixed ordering $1, a, a^{2}, \ldots, a^{2 n-1}, b, a b, \ldots, a^{2 n-1} b, b^{2}, a b^{2}, \ldots, a^{2 n-1} b^{2}, b^{3}, a b^{3}, \ldots, a^{2 n-1} b^{3}$ for the elements of the group $V_{8 n}$. The adjacency matrix of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the one-dimensional irreducible representations $\chi_{1}, \chi_{2}, \chi_{3}$ and $\chi_{4}$.

$$
\begin{aligned}
& u_{1}=\frac{1}{\sqrt{8 n}}(1, \ldots, 1)^{t} \\
& u_{2}=\frac{1}{\sqrt{8 n}}(1, \ldots, 1,-1, \ldots,-1,1, \ldots, 1,-1, \ldots,-1)^{t} \\
& u_{3}=\frac{1}{\sqrt{8 n}}(1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1)^{t} \text { and } \\
& u_{4}=\frac{1}{\sqrt{8 n}}(1,-1, \ldots, 1,-1,-1,1, \ldots,-1,1,1,-1, \ldots, 1,-1,-1,1, \ldots,-1,1)^{t}
\end{aligned}
$$

Now we compute the eigenvectors of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ corresponding to the twodimensional irreducible representations. The adjacency matrix of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the two-dimensional representations $\psi_{j}$, for $0 \leq j \leq n-1$.

$$
\begin{aligned}
& u_{j}^{(1)}=\frac{1}{\sqrt{4 n}}\left(\left\{\omega^{2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\omega^{2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t} \\
& u_{j}^{(2)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{\omega^{2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\omega^{2 r j}\right\}_{r=0}^{2 n-1}\right)^{t} \\
& u_{j}^{(3)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{(-1)^{r+1} \omega^{-2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{(-1)^{r} \omega^{-2 r j}\right\}_{r=0}^{2 n-1}\right)^{t} \text { and } \\
& u_{j}^{(4)}=\frac{1}{\sqrt{4 n}}\left(\left\{(-1)^{r} \omega^{2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{(-1)^{r+1} \omega^{2 r j}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t}
\end{aligned}
$$

The adjacency matrix of the normal Cayley graph $\operatorname{Cay}\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the two-dimensional representations $\phi_{k}$, for $1 \leq k \leq n-1$.

$$
\begin{aligned}
v_{k}^{(1)} & =\frac{1}{\sqrt{4 n}}\left(\left\{\omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t} \\
v_{k}^{(2)} & =\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{\omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{r k}\right\}_{r=0}^{2 n-1}\right)^{t} \\
v_{k}^{(3)} & =\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{\omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{-r k}\right\}_{r=0}^{2 n-1}\right)^{t} \text { and } \\
v_{k}^{(4)} & =\frac{1}{\sqrt{4 n}}\left(\left\{\omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t}
\end{aligned}
$$

### 2.2 Spectra of normal Cayley graphs $\operatorname{Cay}\left(V_{8 n}, S\right)$, for even $n$

The adjacency matrix of the normal Cayley graph $\operatorname{Cay}\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the one-dimensional irreducible representations $\chi_{1}, \ldots, \chi_{8}$.

$$
\begin{aligned}
& u_{1}=\frac{1}{\sqrt{8 n}}(1, \ldots, 1)^{t}, \\
& u_{2}=\frac{1}{\sqrt{4 n}}\left(\left\{\mathbf{i}^{r}\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(-\mathbf{i})\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(-1)\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(\mathbf{i})\right\}_{r=0}^{2 n-1}\right)^{t}, \\
& u_{3}=\frac{1}{\sqrt{8 n}}(1,-1, \ldots, 1,-1,-1,1, \ldots,-1,1,1,-1, \ldots, 1,-1,-1,1, \ldots,-1,1)^{t}, \\
& u_{4}=\frac{1}{\sqrt{4 n}}\left(\left\{(-\mathbf{i})^{r}\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(\mathbf{i})\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(-1)\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(-\mathbf{i})\right\}_{r=0}^{2 n-1}\right)^{t}, \\
& u_{5}=\frac{1}{\sqrt{8 n}}(1, \ldots, 1,-1, \ldots,-1,1, \ldots, 1,-1, \ldots,-1)^{t}, \\
& u_{6}=\frac{1}{\sqrt{4 n}}\left(\left\{\mathbf{i}^{r}\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(\mathbf{i})\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(-1)\right\}_{r=0}^{2 n-1},\left\{\mathbf{i}^{r}(-\mathbf{i})\right\}_{r=0}^{2 n-1}\right)^{t}, \\
& u_{7}=\frac{1}{\sqrt{8 n}}(1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1,1,-1, \ldots, 1,-1)^{t} \text { and } \\
& u_{8}=\frac{1}{\sqrt{4 n}}\left(\left\{(-\mathbf{i})^{r}\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(-\mathbf{i})\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(-1)\right\}_{r=0}^{2 n-1},\left\{(-\mathbf{i})^{r}(\mathbf{i})\right\}_{r=0}^{2 n-1}\right)^{t} .
\end{aligned}
$$

Now we compute the eigenvectors of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ corresponding to the twodimensional irreducible representations. The adjacency matrix of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the two-dimensional representations $\psi_{j}$, for $1 \leq j \leq n-1$.

$$
\begin{aligned}
& u_{j}^{(1)}=\frac{1}{\sqrt{4 n}}\left(\left\{\omega^{r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{r j}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t} \\
& u_{j}^{(2)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{\mathbf{i} \omega^{r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\mathbf{i} \omega^{r j}\right\}_{r=0}^{2 n-1}\right)^{t} \\
& u_{j}^{(3)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{-\mathbf{i} \omega^{-r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\mathbf{i} \omega^{-r j}\right\}_{r=0}^{2 n-1}\right)^{t} \text { and } \\
& u_{j}^{(4)}=\frac{1}{\sqrt{4 n}}\left(\left\{\omega^{-r j}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\omega^{-r j}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t}
\end{aligned}
$$

The adjacency matrix of the normal Cayley graph $\operatorname{Cay}\left(V_{8 n}, S\right)$ has the following eigenvectors corresponding to the two-dimensional representations $\phi_{k}$, for $1 \leq k \leq n-1$.

$$
\begin{aligned}
& v_{k}^{(1)}=\frac{1}{\sqrt{4 n}}\left(\left\{\mathbf{i}^{r} \omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\mathbf{i}^{r} \omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t} \\
& v_{k}^{(2)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{\mathbf{i}^{r} \omega^{r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\mathbf{i}^{r} \omega^{r k}\right\}_{r=0}^{2 n-1}\right)^{t} \\
& v_{k}^{(3)}=\frac{1}{\sqrt{4 n}}\left(\mathbf{0},\left\{-\mathbf{i}^{r} \omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{\mathbf{i}^{r} \omega^{-r k}\right\}_{r=0}^{2 n-1}\right)^{t} \text { and } \\
& v_{k}^{(4)}=\frac{1}{\sqrt{4 n}}\left(\left\{\mathbf{i}^{r} \omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0},\left\{-\mathbf{i}^{r} \omega^{-r k}\right\}_{r=0}^{2 n-1}, \mathbf{0}\right)^{t}
\end{aligned}
$$

## 3 PST on Cayley graphs

Let $\Gamma$ be a simple graph with $n$ vertices and $\operatorname{Spec}(\Gamma)$ denotes the set of all the eigenvalues of $\Gamma$. Let $A$ be the adjacency matrix of $\Gamma$ and $\lambda_{1}, \ldots, \lambda_{n}$ are the eigenvalues of $A$. Let $P=\left(v_{1}, \ldots, v_{n}\right)$ be an unitary matrix, where $v_{i}$ is an eigenvector corresponding to the eigenvalue $\lambda_{i}(1 \leq i \leq n)$. Then the spectral decomposition of $A$ is given by

$$
A=\lambda_{1} E_{1}+\cdots+\lambda_{n} E_{n}
$$

where $E_{i}=v_{i} v_{i}^{*}(1 \leq i \leq n)$ satisfies

$$
E_{i} E_{j}=\left\{\begin{array}{cl}
E_{i} & \text { if } i=j \\
0 & \text { otherwise }
\end{array}\right.
$$

Therefore, the spectral decomposition of the transition matrix $H(t)$ is given by

$$
H(t)=\exp \left(-\mathbf{i} \lambda_{1} t\right) E_{1}+\cdots+\exp \left(-\mathbf{i} \lambda_{n} t\right) E_{n}
$$

The 2-adic exponential valuation of rational numbers is denoted by $\Upsilon_{2}$ and is a mapping defined by

$$
\Upsilon_{2}: \mathbb{Q} \rightarrow \mathbb{Z} \cup\{\infty\}, \text { such that } \Upsilon_{2}(0)=\infty, \text { and } \Upsilon_{2}\left(2^{l} \frac{a}{b}\right)=l, \text { where } a, b, l \in \mathbb{Z} \text { and } 2 \nmid a b .
$$

We assume that $\infty+\infty=\infty+l=\infty$ and $\infty>l$ for any $l \in \mathbb{Z}$. Then $\Upsilon_{2}$ has the following properties. For $\beta, \beta^{\prime} \in \mathbb{Q}$,

1. $\Upsilon_{2}\left(\beta \beta^{\prime}\right)=\Upsilon_{2}(\beta)+\Upsilon_{2}\left(\beta^{\prime}\right)$ and
2. $\Upsilon_{2}\left(\beta+\beta^{\prime}\right) \geq \min \left(\Upsilon_{2}(\beta), \Upsilon_{2}\left(\beta^{\prime}\right)\right)$ and the equality holds if $\Upsilon_{2}(\beta) \neq \Upsilon_{2}\left(\beta^{\prime}\right)$.

We write the vertex set of $\operatorname{Cay}\left(V_{8 n}, S\right)$ as $V_{1} \cup V_{2} \cup V_{3} \cup V_{4}$, where

$$
\begin{aligned}
& V_{1}=\{0,1, \ldots 2 n-1\}, \\
& V_{2}=\{2 n, 2 n+1, \ldots, 4 n-1\}, \\
& V_{3}=\{4 n, 4 n+1, \ldots, 6 n-1\} \text { and } \\
& V_{4}=\{6 n, 6 n+1, \ldots, 8 n-1\} .
\end{aligned}
$$

### 3.1 PST on normal Cayley graphs over the group $V_{8 n}$, for odd $n$

We want to state the main result of this section.

Theorem 3.1. Let $S$ be a non-empty subset of $V_{8 n}$ such that $1 \notin S$ and $S g=g S$ for all $g \in V_{8 n}$. Let $\Gamma=\operatorname{Cay}\left(V_{8 n}, S\right)$ be a connected Cayley graph with connection set $S$, where $n$ is odd. Then $\Gamma$ has four distinct eigenvalues which corresponds to the one-dimensional representations $\chi_{1}, \chi_{2}, \chi_{3}$ and $\chi_{4}$, respectively, with one is $\alpha_{1}=|S|$ and the other three eigenvalues are denoted by $\alpha_{2}, \alpha_{3}$ and $\alpha_{4}$, and some
multiple eigenvalues corresponding to the two-dimensional representations $\psi_{j}$ and $\phi_{k}$, denoted by $\beta_{j}$ and $\gamma_{k}$, respectively, for $0 \leq j \leq n-1$ and $1 \leq k \leq n-1$.

1. If $u \in V_{1}, v \in V_{2}$ or $u \in V_{1}, v \in V_{4}$ or $u \in V_{2}, v \in V_{1}$ or $u \in V_{2}, v \in V_{3}$ or $u \in V_{3}, v \in V_{2}$ or $u \in V_{3}, v \in V_{4}$ or $u \in V_{4}, v \in V_{1}$ or $u \in V_{4}, v \in V_{3}$, then $\Gamma$ cannot have PST between two distinct vertices $u$ and $v$.
2. If $u, v \in V_{1}$ or $u, v \in V_{2}$ or $u, v \in V_{3}$ or $u, v \in V_{4}$, then $\Gamma$ cannot have PST between two distinct vertices $u$ and $v$.
3. If $u \in V_{1}, v \in V_{3}$ or $u \in V_{2}, v \in V_{4}$ or $u \in V_{3}, v \in V_{1}$ or $u \in V_{4}, v \in V_{2}$, then $\Gamma$ has PST between the vertices $u$ and $v$ if and only if the following three conditions hold.
(i) All the eigenvalues of $\Gamma$ are integers, namely, $\Gamma$ is integral,
(ii) $u=v+4 n$ and
(iii) $\Upsilon_{2}\left(\alpha_{1}-\beta_{j}\right)$ is a constant, say $\mu$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{k}\right)$ are all bigger than $\mu$, for $0 \leq j \leq n-1$ and $1 \leq k \leq n-1$.

Furthermore, when the conditions (i), (ii) and (iii) hold, the minimum time at which $\Gamma$ has PST between $u$ and $v$ is $\frac{\pi}{M}$, where $M=\operatorname{gcd}\left(\alpha-\alpha_{1}: \alpha \in \operatorname{Spec}(\Gamma) \backslash\left\{\alpha_{1}\right\}\right)$.

Proof. The adjacency matrix $A$ of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ has the eigenvectors $u_{i}, u_{j}{ }^{(i)}$ and $v_{k}{ }^{(i)}(1 \leq i \leq 4,0 \leq j \leq n-1$ and $1 \leq k \leq n-1)$ which are introduced in section 2.1. Hence we have the following unitary matrix

$$
\begin{array}{r}
P=\left(u_{1}, u_{2}, u_{3}, u_{4}, u_{0}^{(1)}, u_{0}^{(2)}, u_{0}^{(3)}, u_{0}^{(4)}, \ldots, u_{n-1}^{(1)}, u_{n-1}^{(2)}, u_{n-1}^{(3)}, u_{n-1}^{(4)},\right. \\
\left.v_{1}^{(1)}, v_{1}^{(2)}, v_{1}^{(3)}, v_{1}^{(4)}, \ldots, v_{n-1}^{(1)}, v_{n-1}^{(2)}, v_{n-1}(3), v_{n-1}^{(4)}\right) .
\end{array}
$$

Let $J_{m}$ be the all-one matrix of order $m$. Then the projective matrices corresponding to the eigenvectors $u_{1}, u_{2}, u_{3}$ and $u_{4}$ are given by

$$
\begin{aligned}
& E_{1}= u_{1} u_{1}^{*}= \\
& E_{2}= \frac{1}{8 n} J_{8 n}, \\
&\left(\begin{array}{cccc}
J_{2 n} & -J_{2 n} & J_{2 n} & -J_{2 n} \\
-J_{2 n} & J_{2 n} & -J_{2 n} & J_{2 n} \\
J_{2 n} & -J_{2 n} & J_{2 n} & -J_{2 n} \\
-J_{2 n} & J_{2 n} & -J_{2 n} & J_{2 n}
\end{array}\right) \\
& E_{3}=u_{3} u_{3}^{*}=\frac{1}{8 n}\left((-1)^{u+v}\right), u, v \in\{0,1 \ldots, 8 n-1\} \text { and } \\
& E_{4}=u_{4} u_{4}^{*}=\frac{1}{8 n}\left(e_{4}(u, v)\right),
\end{aligned}
$$

where
(i) $e_{4}(u, v)=(-1)^{u+v}$, when $u, v \in V_{1} \cup V_{3}$ or $u, v \in V_{2} \cup V_{4}$
(ii) $e_{4}(u, v)=(-1)^{u+v+1}$, when $u \in V_{1} \cup V_{3}, v \in V_{2} \cup V_{4}$ or $u \in V_{2} \cup V_{4}, v \in V_{1} \cup V_{3}$,

The projective matrices corresponding to the eigenvectors $u_{j}^{(1)}, u_{j}^{(2)}, u_{j}^{(3)}$ and $u_{j}^{(4)}$, where $0 \leq j \leq n-1$ are given by

$$
\begin{aligned}
& E_{j}^{(1)}=u_{j}^{(1)} u_{j}^{(1)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
X_{1} & 0 & -X_{1} & 0 \\
0 & 0 & 0 & 0 \\
-X_{1} & 0 & X_{1} & 0 \\
0 & 0 & 0 & 0
\end{array}\right), E_{j}^{(2)}=u_{j}^{(2)} u_{j}^{(2)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & X_{1} & 0 & -X_{1} \\
0 & 0 & 0 & 0 \\
0 & -X_{1} & 0 & X_{1}
\end{array}\right), \\
& E_{j}^{(3)}=u_{j}^{(3)} u_{j}^{(3)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & X_{2} & 0 & -X_{2} \\
0 & 0 & 0 & 0 \\
0 & -X_{2} & 0 & X_{2}
\end{array}\right) \text { and } E_{j}^{(4)}=u_{j}^{(4)} u_{j}^{(4)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
X_{2} & 0 & -X_{2} & 0 \\
0 & 0 & 0 & 0 \\
-X_{2} & 0 & X_{2} & 0 \\
0 & 0 & 0 & 0
\end{array}\right) .
\end{aligned}
$$

The projective matrices corresponding to the eigenvectors $v_{k}^{(1)}, v_{k}^{(2)}, v_{k}^{(3)}$ and $v_{k}^{(4)}$, where $1 \leq k \leq n-1$ are given by

$$
\begin{aligned}
& F_{k}^{(1)}=v_{k}^{(1)} v_{k}^{(1)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
Y_{1} & 0 & Y_{1} & 0 \\
0 & 0 & 0 & 0 \\
Y_{1} & 0 & Y_{1} & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad F_{k}^{(2)}=v_{k}^{(2)} v_{k}^{(2)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & Y_{1} & 0 & Y_{1} \\
0 & 0 & 0 & 0 \\
0 & Y_{1} & 0 & Y_{1}
\end{array}\right), \\
& F_{k}^{(3)}=v_{k}^{(3)} v_{k}^{(3)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & Y_{2} & 0 & Y_{2} \\
0 & 0 & 0 & 0 \\
0 & Y_{2} & 0 & Y_{2}
\end{array}\right) \quad \text { and } \quad F_{k}^{(4)}=v_{k}^{(4)} v_{k}^{(4)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
Y_{2} & 0 & Y_{2} & 0 \\
0 & 0 & 0 & 0 \\
Y_{2} & 0 & Y_{2} & 0 \\
0 & 0 & 0 & 0
\end{array}\right),
\end{aligned}
$$

where

$$
\left.\begin{array}{l}
X_{1}=\left(\begin{array}{cccc}
1 & \omega^{-2 j} & \cdots & \omega^{-2(2 n-1) j} \\
\omega^{2 j} & 1 & \cdots & \omega^{-2(2 n-2) j} \\
\vdots & \vdots & \vdots & \vdots \\
\omega^{2(2 n-1) j} & \omega^{2(2 n-2) j} & \cdots & 1
\end{array}\right), \quad X_{2}=\left(\begin{array}{ccc}
-\omega^{2 j} & \cdots & -\omega^{2(2 n-1) j} \\
1 & \cdots & \omega^{2(2 n-2) j} \\
-\omega^{-2 j} & \vdots & \vdots \\
\vdots & \vdots \\
-\omega^{-2(2 n-1) j} & \omega^{-2(2 n-2) j} & \cdots
\end{array}\right] 1
\end{array}\right), ~ \begin{array}{cccc}
1 & \omega^{-(2 n-1) k} \\
\left.Y_{1}=\left(\begin{array}{ccc}
\omega^{-k} & \cdots & \\
\omega^{k} & 1 & \cdots \\
\vdots & \vdots & \vdots \\
\omega^{-(2 n-2) k} \\
\omega^{(2 n-1) k} & \omega^{(2 n-2) k} & \cdots
\end{array}\right) \quad \begin{array}{c}
1 \\
1
\end{array}\right)
\end{array}
$$

Therefore, the transition matrix $H(t)$ of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ is given by

$$
\begin{aligned}
H(t) & =\exp \left(-\mathbf{i} t \alpha_{1}\right) E_{1}+\exp \left(-\mathbf{i} t \alpha_{2}\right) E_{2}+\exp \left(-\mathbf{i} t \alpha_{3}\right) E_{3}+\exp \left(-\mathbf{i} t \alpha_{4}\right) E_{4} \\
& +\sum_{j=0}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(E_{j}^{(1)}+E_{j}^{(2)}+E_{j}^{(3)}+E_{j}^{(4)}\right)+\sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(F_{k}^{(1)}+F_{k}^{(2)}+F_{k}^{(3)}+F_{k}^{(4)}\right)
\end{aligned}
$$

Now we compute the $(u, v)$-th element of the transition matrix. We have the following three cases.
Case 1. If $u \in V_{1}, v \in V_{2}$ or $u \in V_{1}, v \in V_{4}$ or $u \in V_{2}, v \in V_{1}$ or $u \in V_{2}, v \in V_{3}$ or $u \in V_{3}, v \in V_{2}$ or $u \in V_{3}, v \in V_{4}$ or $u \in V_{4}, v \in V_{1}$ or $u \in V_{4}, v \in V_{3}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)-\exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right)+(-1)^{u+v+1} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right) \\
& +\frac{1}{4 n} \sum_{j=0}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)(0+0+0+0)+\frac{1}{4 n} \sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)(0+0+0+0) \\
& =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)-\exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{4}\right)+(-1)^{u+v+1} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right) . \\
& \text { This implies that }\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 4=\frac{1}{2 n}<1 .
\end{aligned}
$$

Therefore, PST cannot occur in this case.
Case 2. If $u, v \in V_{1}$ or $u, v \in V_{2}$ or $u, v \in V_{3}$ or $u, v \in V_{4}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)+\exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right) \\
& +\frac{1}{4 n} \sum_{j=0}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(\omega^{2(u-v) j}+(-1)^{u+v} \omega^{2(v-u) j}\right)+\frac{1}{4 n} \sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(\omega^{(u-v) k}+\omega^{(v-u) k}\right)
\end{aligned}
$$

This implies that $\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 4+\frac{1}{4 n} \times 2 n+\frac{1}{4 n} \times 2(n-1)=\frac{2}{4 n}+\frac{2 n+2 n-2}{4 n}=\frac{4 n}{4 n}=1$
Therefore, $\left|H(t)_{u v}\right| \leq 1$. Thus $\left|H(t)_{u v}\right|=1$ if and only if for $0 \leq j \leq n-1$ and $1 \leq k \leq n-1$, it holds that,

$$
\begin{aligned}
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\exp \left(-\mathbf{i} t \alpha_{2}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{4}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{2(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \omega^{2(v-u) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(v-u) k} \exp \left(-\mathbf{i} t \gamma_{k}\right)
\end{aligned}
$$

From the last two equations, since $\omega$ is the $2 n$-th root of unity, we get that $n$ divides $u-v$. Without loss of generality, we assume that $u \geq v$. This implies that either $u=v$ or $u-v=n$. Let $t=2 \pi T$. Then we have

$$
\begin{align*}
& \left(\alpha_{1}-\alpha_{2}\right) T \in \mathbb{Z}  \tag{1}\\
& \left(\alpha_{1}-\alpha_{3}\right) T-\frac{u+v}{2} \in \mathbb{Z}  \tag{2}\\
& \left(\alpha_{1}-\alpha_{4}\right) T-\frac{u+v}{2} \in \mathbb{Z}  \tag{3}\\
& \left(\alpha_{1}-\beta_{j}\right) T \in \mathbb{Z}  \tag{4}\\
& \left(\alpha_{1}-\beta_{j}\right) T-\frac{u+v}{2} \in \mathbb{Z}  \tag{5}\\
& \left(\alpha_{1}-\gamma_{k}\right) T-\frac{k}{2} \in \mathbb{Z}  \tag{6}\\
& \left(\alpha_{1}-\gamma_{k}\right) T+\frac{k}{2} \in \mathbb{Z} \tag{7}
\end{align*}
$$

Since $0=\operatorname{tr}(A)=\alpha_{1}+\alpha_{2}+\alpha_{3}+\alpha_{4}+4 \sum_{j=0}^{n-1} \beta_{j}+4 \sum_{k=1}^{n-1} \gamma_{k}$, we have that $8 n \alpha_{1} T \in \mathbb{Z}$, and since $\alpha_{1}=|S|$ is a positive integer, we have $T \in \mathbb{Q}$. This implies that all the eigenvalues of the graph are rational numbers. It is well known that any rational eigenvalue of a graph is an integer. Therefore in this case the graph is integral.

From (4) and (5) it follows that $u+v$ is even. Since $u+v$ and $u-v$ have the same parity, so $u-v$ cannot be equal to $n$. Hence, $u=v$. Therefore, $\Gamma$ cannot have PST between distinct vertices $u$ and $v$.
Case 3. If $u \in V_{1}, v \in V_{3}$ or $u \in V_{2}, v \in V_{4}$ or $u \in V_{3}, v \in V_{1}$ or $u \in V_{4}, v \in V_{2}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)+\exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right)+\exp \left(-\mathbf{i} t \alpha_{4}\right)\right) \\
& +\frac{1}{4 n} \sum_{j=0}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(-\omega^{2(u-v) j}+(-1)^{u+v+1} \omega^{2(v-u) j}\right)+\frac{1}{4 n} \sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(\omega^{(u-v) k}+\omega^{(v-u) k}\right)
\end{aligned}
$$

This implies that $\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 4+\frac{1}{4 n} \times 2 n+\frac{1}{4 n} \times 2(n-1)=\frac{2}{4 n}+\frac{2 n+2 n-2}{4 n}=\frac{4 n}{4 n}=1$. Therefore, $\left|H(t)_{u v}\right| \leq 1$. Thus, $\left|H(t)_{u v}\right|=1$ if and only if for $0 \leq j \leq n-1$ and $1 \leq k \leq n-1$, it holds that,

$$
\begin{aligned}
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\exp \left(-\mathbf{i} t \alpha_{2}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{4}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-\omega^{2(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v+1} \omega^{2(v-u) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(v-u) k} \exp \left(-\mathbf{i} t \gamma_{k}\right)
\end{aligned}
$$

From the last two equations, since $\omega$ is the $2 n$-th root of unity, we get that $n$ divides $u-v$. Without loss of generality, we assume that $u>v$. This implies that either $u-v=3 n$ or $u-v=4 n$ or $u-v=5 n$. Let $t=2 \pi T$,. Then we have

$$
\begin{align*}
& \left(\alpha_{1}-\alpha_{2}\right) T \in \mathbb{Z}  \tag{8}\\
& \left(\alpha_{1}-\alpha_{3}\right) T-\frac{u+v}{2} \in \mathbb{Z}  \tag{9}\\
& \left(\alpha_{1}-\alpha_{4}\right) T-\frac{u+v}{2} \in \mathbb{Z}  \tag{10}\\
& \left(\alpha_{1}-\beta_{j}\right) T+\frac{1}{2} \in \mathbb{Z}  \tag{11}\\
& \left(\alpha_{1}-\beta_{j}\right) T+\frac{u+v+1}{2} \in \mathbb{Z}  \tag{12}\\
& \left(\alpha_{1}-\gamma_{k}\right) T+\frac{(u-v) k}{2 n} \in \mathbb{Z}  \tag{13}\\
& \left(\alpha_{1}-\gamma_{k}\right) T+\frac{(v-u) k}{2 n} \in \mathbb{Z} \tag{14}
\end{align*}
$$

Since $0=\operatorname{tr}(A)=\alpha_{1}+\alpha_{2}+\alpha_{3}+\alpha_{4}+4 \sum_{j=0}^{n-1} \beta_{j}+4 \sum_{k=1}^{n-1} \gamma_{k}$, we have that $8 n \alpha_{1} T \in \mathbb{Z}$, and since $\alpha_{1}=|S|$ is a positive integer, we have $T \in \mathbb{Q}$. This implies that all the eigenvalues of the graph are rational numbers. It is well known that any rational eigenvalue of a graph is an integer. Therefore in this case the graph is integral.

From (11) and (12) it follows that $u+v$ is even. Since $u+v$ and $u-v$ have the same parity, so $u-v$ is neither $3 n$ nor $5 n$. Therefore $u-v=4 n$.
Conditions (8) to (14) can be written as follows

$$
\begin{align*}
& \left(\alpha_{1}-\alpha_{2}\right) T \in \mathbb{Z}  \tag{15}\\
& \left(\alpha_{1}-\alpha_{3}\right) T \in \mathbb{Z}  \tag{16}\\
& \left(\alpha_{1}-\alpha_{4}\right) T \in \mathbb{Z}  \tag{17}\\
& \left(\alpha_{1}-\beta_{j}\right) T-\frac{1}{2} \in \mathbb{Z}  \tag{18}\\
& \left(\alpha_{1}-\gamma_{k}\right) T \in \mathbb{Z} \tag{19}
\end{align*}
$$

Suppose that $\left(\alpha_{1}-\beta_{r}\right) T,\left(\alpha_{1}-\beta_{s}\right) T \in \frac{1}{2}+\mathbb{Z}$, for $r, s \in\{0, \ldots, n-1\}$, then $\Upsilon_{2}\left(\left(\alpha_{1}-\beta_{r}\right) T\right)=$ $\Upsilon_{2}\left(\left(\alpha_{1}-\beta_{s}\right) T\right)=-1$. This implies that $\Upsilon_{2}\left(\alpha_{1}-\beta_{r}\right)=\Upsilon_{2}\left(\alpha_{1}-\beta_{s}\right)=-1-\Upsilon_{2}(T)$. Thus for all $j \in\{0, \ldots, n-1\}, \Upsilon_{2}\left(\alpha_{1}-\beta_{j}\right)$ is a constant, $\mu=-1-\Upsilon_{2}(T)$. From (15) it follows that $\Upsilon_{2}\left(\left(\alpha_{1}-\alpha_{2}\right) T\right) \geq 0$. Therefore, $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right) \geq \mu+1$. Simillarly it can be shown that $\Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{k}\right)$ are also bigger than $\mu$, for all $k \in\{1, \ldots, n-1\}$.

Suppose that $(i),(i i)$ and (iii) hold. Let $M_{1}=\operatorname{gcd}\left(\alpha_{1}-\alpha_{2}, \alpha_{1}-\alpha_{3}, \alpha_{1}-\alpha_{4}, \alpha_{1}-\gamma_{k}: 1 \leq k \leq n-1\right)$ and $M_{2}=\operatorname{gcd}\left(\alpha_{1}-\beta_{j}: 0 \leq j \leq n-1\right)$. It can be easily seen that $\Upsilon_{2}\left(M_{2}\right)=\mu$.

Then the conditions (15), (16), (17) and (19) implies that $T \in \frac{1}{M_{1}} \mathbb{Z}$ and condition (18) imply that
$T \in \frac{1}{M_{2}}\left(\frac{1}{2}+\mathbb{Z}\right)$. let $\tilde{T}=\{t>0: \Gamma$ has PST between $u$ and $v$ at time $t\}$. Then we get

$$
\begin{align*}
\tilde{T} & =\left(\frac{2 \pi}{M_{1}} \mathbb{Z}\right) \cap\left(\frac{2 \pi}{M_{2}}\left(\frac{1}{2}+\mathbb{Z}\right)\right) \cap \mathbb{R}_{>0}  \tag{20}\\
& =\frac{\pi}{M_{1} M_{2}}\left(2 M_{2} \mathbb{Z} \cap M_{1}(1+2 \mathbb{Z})\right) \cap \mathbb{R}_{>0} \tag{21}
\end{align*}
$$

For every $z \in \mathbb{Z}$, it is easy to check that

$$
\begin{aligned}
& z \in 2 M_{2} \mathbb{Z} \cap M_{1}(1+2 \mathbb{Z}) \\
& \Leftrightarrow z=2 M_{2} x_{0}=M_{1}\left(1+2 y_{0}\right), \text { for some } x_{0}, y_{0} \in \mathbb{Z} \\
& \Leftrightarrow 2 M_{2} x-2 M_{1} y=M_{1} \text { has a solution } \\
& \Leftrightarrow \operatorname{gcd}\left(2 M_{1}, 2 M_{2}\right) \mid M_{1} \\
& \Leftrightarrow \Upsilon_{2}\left(M_{1}\right) \geq \mu+1\left(\text { since } \Upsilon_{2}\left(M_{2}\right)=\mu\right)
\end{aligned}
$$

Let $M=\operatorname{gcd}\left(M_{1}, M_{2}\right)$. Write $M_{1}=m_{1} M, M_{2}=m_{2} M$. Then $\operatorname{gcd}\left(m_{1}, m_{2}\right)=1$. From $\Upsilon_{2}\left(M_{1}\right) \geq$ $\mu+1, \Upsilon_{2}\left(M_{2}\right)=\mu$, we get that $\Upsilon_{2}(M)=\mu$ and $m_{1}$ is even, $m_{2}$ is odd. Then $2 M_{2} x-2 M_{1} y=M_{1} \Leftrightarrow$ $m_{2} x-m_{1} y=\frac{m_{1}}{2}$. Since $\operatorname{gcd}\left(m_{1}, m_{2}\right)=1$, so the solutions of the Diophantine equation $m_{2} x-m_{1} y=\frac{m_{1}}{2}$ are given by

$$
\begin{aligned}
& x=\frac{m_{1}}{2}+m_{1} l \\
& y=\frac{m_{2}-1}{2}+m_{2} l, \text { where } l \in \mathbb{Z}
\end{aligned}
$$

Thus

$$
z=2 M_{2} x=2 M_{2} \frac{m_{1}}{2}(1+2 l)=\frac{M_{1} M_{2}}{M}(1+2 l)
$$

and

$$
2 M_{2} \mathbb{Z} \cap M_{1}(1+2 \mathbb{Z})=\frac{M_{1} M_{2}}{M}(1+2 Z)
$$

By (21), we get

$$
\begin{aligned}
\tilde{T} & =\left(\frac{\pi}{M}+\frac{2 \pi}{M} \mathbb{Z}\right) \cap \mathbb{R}_{>0} \\
& =\left\{\frac{\pi}{M}+\frac{2 \pi}{M} l: l \in \mathbb{N} \cup\{0\}\right\}
\end{aligned}
$$

Therefore, the minimum time at which $\Gamma$ has PST between $u$ and $v$ is $\frac{\pi}{M}$.
This completes the proof.

### 3.2 PST on normal Cayley graphs over the group $V_{8 n}, n$ even

The main result of this section is as follows.
Theorem 3.2. Let $S$ be a non-empty subset of $V_{8 n}$ such that $\mathbf{1} \notin S$ and $S g=g S$ for all $g \in V_{8 n}$. Let $\Gamma=\operatorname{Cay}\left(V_{8 n}, S\right)$ be a connected Cayley graph with connection set $S$, where $n$ is even. Then $\Gamma$ has eight distinct eigenvalues which corresponds to the one-dimensional representations $\chi_{1}, \ldots, \chi_{8}$, respectively, with one is $\alpha_{1}=|S|$ and the other three eigenvalues are denoted by $\alpha_{2}, \alpha_{3}, \alpha_{4}, \alpha_{5}, \alpha_{6}, \alpha_{7}$ and $\alpha_{8}$, and some multiple eigenvalues corresponding to the two-dimensional representations $\psi_{j}$ and $\phi_{k}$, denoted by $\beta_{j}$ and $\gamma_{k}$, respectively, for $1 \leq j \leq n-1$ and $1 \leq k \leq n-1$.

1. If $u \in V_{1}, v \in V_{2}$ or $u \in V_{1}, v \in V_{4}$ or $u \in V_{2}, v \in V_{1}$ or $u \in V_{2}, v \in V_{3}$ or $u \in V_{3}, v \in V_{2}$ or $u \in V_{3}, v \in V_{4}$ or $u \in V_{4}, v \in V_{1}$ or $u \in V_{4}, v \in V_{3}$, then $\Gamma$ cannot have PST between distinct vertices $u$ and $v$.
2. If $u, v \in V_{1}$ or $u, v \in V_{2}$ or $u, v \in V_{3}$ or $u, v \in V_{4}$, then $\Gamma$ has PST between distinct vertices $u$ and $v$ if and only if the following three conditions hold.
(i) All the eigenvalues of $\Gamma$ are integers.
(ii) $u=v+n$.
(iii) (a) If $n \equiv 0(\bmod 4)$, then $\Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}-1}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}-1}\right)$ are the same, say $\mu_{1}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right), \Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}}\right)$ are all strictly greater than $\mu_{1}$, for $1 \leq j^{\prime} \leq \frac{n-1}{2}$ and $1 \leq k^{\prime} \leq \frac{n-1}{2}$.
(b) If $n \equiv 2(\bmod 4)$, then $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right), \Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}-1}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}}\right)$ are the same, say $\mu_{2}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right), \Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}-1}\right)$ are all strictly greater than $\mu_{2}$, for $1 \leq j^{\prime} \leq \frac{n-1}{2}$ and $1 \leq k^{\prime} \leq \frac{n-1}{2}$.
3. If $u \in V_{1}, v \in V_{3}$ or $u \in V_{2}, v \in V_{4}$ or $u \in V_{3}, v \in V_{1}$ or $u \in V_{4}, v \in V_{2}$, then $\Gamma$ has PST between distinct vertices $u$ and $v$ if and only if the following three conditions hold.
(i) All the eigenvalues of $\Gamma$ are integers.
(ii) $u=v+4 n$.
(iii) $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{k}\right)$ are the same, say $\mu_{3}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\beta_{j}\right)$ are all strictly greater than $\mu_{3}$, for $1 \leq j \leq n-1$ and $1 \leq k \leq n-1$.

Furthermore, the minimum time at which $\Gamma$ has PST between $u$ and $v$ is $\frac{\pi}{M}$, where $M=\operatorname{gcd}(\alpha-$ $\left.\alpha_{1}: \alpha \in \operatorname{Spec}(\Gamma) \backslash\left\{\alpha_{1}\right\}\right)$.

Proof. The adjacency matrix $A$ of the normal Cayley graph Cay $\left(V_{8 n}, S\right)$ has the eigenvectors $u_{i}, u_{j}{ }^{(i)}$ and $v_{k}{ }^{(i)}(1 \leq i \leq 8,1 \leq j \leq n-1$ and $1 \leq k \leq n-1)$ which are introduced in subsection 2.2. Hence
we have the following unitary matrix

$$
\begin{array}{r}
P=\left(u_{1}, u_{2}, u_{3}, u_{4}, u_{5}, u_{6}, u_{7}, u_{8}, u_{1}^{(1)}, u_{1}^{(2)}, u_{1}^{(3)}, u_{1}^{(4)}, \ldots, u_{n-1}^{(1)}, u_{n-1}{ }^{(2)}, u_{n-1}{ }^{(3)}, u_{n-1}{ }^{(4)},\right. \\
\left.v_{1}^{(1)}, v_{1}^{(2)}, v_{1}^{(3)}, v_{1}^{(4)}, \ldots, v_{n-1}{ }^{(1)}, v_{n-1}{ }^{(2)}, v_{n-1}(3), v_{n-1}{ }^{(4)}\right) .
\end{array}
$$

Let $J_{m}$ be the all-one matrix of order $m$. Then the projective matrices corresponding to the eigenvectors $u_{1}, u_{3}, u_{5}, u_{7}$ are given by

$$
\begin{aligned}
& E_{1}=u_{1} u_{1}^{*}=\frac{1}{8 n} J_{8 n}, \\
& E_{3}=u_{4} u_{4}^{*}=\frac{1}{8 n}\left(e_{3}(u, v)\right), \\
& E_{5}=u_{2} u_{2}^{*}=\frac{1}{8 n}\left(\begin{array}{cccc}
J_{2 n} & -J_{2 n} & J_{2 n} & -J_{2 n} \\
-J_{2 n} & J_{2 n} & -J_{2 n} & J_{2 n} \\
J_{2 n} & -J_{2 n} & J_{2 n} & -J_{2 n} \\
-J_{2 n} & J_{2 n} & -J_{2 n} & J_{2 n}
\end{array}\right), \\
& E_{7}=u_{3} u_{3}^{*}=\frac{1}{8 n}\left((-1)^{u+v}\right), u, v \in\{0,1 \ldots, 8 n-1\}
\end{aligned}
$$

where
(i) $e_{4}(u, v)=(-1)^{u+v}$, when $u, v \in V_{1} \cup V_{3}$ or $u, v \in V_{2} \cup V_{4}$
(ii) $e_{4}(u, v)=(-1)^{u+v+1}$, when $u \in V_{1} \cup V_{3}, v \in V_{2} \cup V_{4}$ or $u \in V_{2} \cup V_{4}, v \in V_{1} \cup V_{3}$,
and the projective matrices corresponding to the eigenvectors $u_{2}, u_{4}, u_{6}$ and $u_{8}$ are given by

$$
\begin{aligned}
& E_{2}=u_{2} u_{2}{ }^{*}=\frac{1}{8 n}\left(\begin{array}{cccc}
X & \mathbf{i} X & -X & -\mathbf{i} X \\
-\mathbf{i} X & X & \mathbf{i} X & -X \\
-X & -\mathbf{i} X & X & \mathbf{i} X \\
\mathbf{i} X & -X & -\mathbf{i} X & X
\end{array}\right), E_{4}=u_{4} u_{4}{ }^{*}=\frac{1}{8 n}\left(\begin{array}{cccc}
Y & -\mathbf{i} Y & -Y & \mathbf{i} Y \\
\mathbf{i} Y & Y & -\mathbf{i} Y & -Y \\
-Y & \mathbf{i} Y & Y & -\mathbf{i} Y \\
-\mathbf{i} Y & -Y & \mathbf{i} Y & Y
\end{array}\right), \\
& E_{6}=u_{6} u_{6}{ }^{*}=\frac{1}{8 n}\left(\begin{array}{cccc}
X & -\mathbf{i} X & -X & \mathbf{i} X \\
\mathbf{i} X & X & -\mathbf{i} X & -X \\
-X & \mathbf{i} X & X & -\mathbf{i} X \\
-\mathbf{i} X & -X & \mathbf{i} X & X
\end{array}\right) \text { and } E_{8}=u_{8} u_{8}^{*}=\frac{1}{8 n}\left(\begin{array}{cccc}
Y & \mathbf{i} Y & -Y & -\mathbf{i} Y \\
-\mathbf{i} Y & Y & \mathbf{i} Y & -Y \\
-Y & -\mathbf{i} Y & Y & \mathbf{i} Y \\
\mathbf{i} Y & -Y & -\mathbf{i} Y & Y
\end{array}\right),
\end{aligned}
$$

where $X_{u v}=\mathbf{i}^{u-v}$ and $Y_{u v}=(-\mathbf{i})^{u-v}$.
The projective matrices corresponding to the eigenvectors $u_{j}^{(1)}, u_{j}^{(2)}, u_{j}^{(3)}$ and $u_{j}^{(4)}$, where $1 \leq j \leq n-1$ are given by

$$
\begin{aligned}
& E_{j}^{(1)}=u_{j}^{(1)} u_{j}^{(1)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
X_{1} & 0 & X_{1} & 0 \\
0 & 0 & 0 & 0 \\
X_{1} & 0 & X_{1} & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad E_{j}^{(2)}=u_{j}^{(2)} u_{j}^{(2)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & X_{1} & 0 & X_{1} \\
0 & 0 & 0 & 0 \\
0 & X_{1} & 0 & X_{1}
\end{array}\right), \\
& E_{j}^{(3)}=u_{j}^{(3)} u_{j}^{(3)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & X_{2} & 0 & X_{2} \\
0 & 0 & 0 & 0 \\
0 & X_{2} & 0 & X_{2}
\end{array}\right) \text { and } E_{j}^{(4)}=u_{j}^{(4)} u_{j}^{(4)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
X_{2} & 0 & X_{2} & 0 \\
0 & 0 & 0 & 0 \\
X_{2} & 0 & X_{2} & 0 \\
0 & 0 & 0 & 0
\end{array}\right) .
\end{aligned}
$$

The projective matrices corresponding to the eigenvectors $v_{k}^{(1)}, v_{k}^{(2)}, v_{k}^{(3)}$ and $v_{k}^{(4)}$, where $1 \leq k \leq n-1$ are given by

$$
\begin{gathered}
F_{k}^{(1)}=v_{k}^{(1)} v_{k}^{(1)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
Y_{1} & 0 & -Y_{1} & 0 \\
0 & 0 & 0 & 0 \\
-Y_{1} & 0 & Y_{1} & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \quad F_{k}^{(2)}=v_{k}^{(2)} v_{k}^{(2)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & Y_{1} & 0 & -Y_{1} \\
0 & 0 & 0 & 0 \\
0 & -Y_{1} & 0 & Y_{1}
\end{array}\right), \\
F_{k}^{(3)}=v_{k}^{(3)} v_{k}^{(3)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & Y_{2} & 0 & -Y_{2} \\
0 & 0 & 0 & 0 \\
0 & -Y_{2} & 0 & Y_{2}
\end{array}\right) \quad \text { and } \quad F_{k}{ }^{(4)}=v_{k}^{(4)} v_{k}^{(4)^{*}}=\frac{1}{4 n}\left(\begin{array}{cccc}
Y_{2} & 0 & -Y_{2} & 0 \\
0 & 0 & 0 & 0 \\
-Y_{2} & 0 & Y_{2} & 0 \\
0 & 0 & 0 & 0
\end{array}\right),
\end{gathered}
$$

where

$$
X_{1}=\left(\begin{array}{cccc}
1 & \omega^{-j} & \cdots & \omega^{-(2 n-1) j} \\
\omega^{j} & 1 & \cdots & \omega^{-(2 n-2) j} \\
\vdots & \vdots & \vdots & \vdots \\
\omega^{(2 n-1) j} & \omega^{(2 n-2) j} & \cdots & 1
\end{array}\right), \quad X_{2}=\left(\begin{array}{cccc}
1 & \omega^{j} & \cdots & \omega^{(2 n-1) j} \\
\omega^{-j} & 1 & \cdots & \omega^{(2 n-2) j} \\
\vdots & \vdots & \vdots & \vdots \\
\omega^{-(2 n-1) j} & \omega^{-(2 n-2) j} & \cdots & 1
\end{array}\right),
$$

$$
\begin{aligned}
& Y_{1}=\left(\begin{array}{cccc}
1 & \mathbf{i}^{-1} \omega^{-k} & \cdots & \mathbf{i}^{-(2 n-1)} \omega^{-(2 n-1) k} \\
\mathbf{i} \omega^{k} & 1 & \ldots & \mathbf{i}^{-(2 n-2)} \omega^{-(2 n-2) k} \\
\vdots & \vdots & \vdots & \vdots \\
\mathbf{i}^{(2 n-1)} \omega^{(2 n-1) k} & \mathbf{i}^{(2 n-2)} \omega^{(2 n-2) k} & \ldots & 1
\end{array}\right) \\
& Y_{2}=\left(\begin{array}{cccc}
1 & \mathbf{i}^{-1} \omega^{k} & \ldots & \mathbf{i}^{-(2 n-1)} \omega^{(2 n-1) k} \\
\mathbf{i} \omega^{-k} & 1 & \ldots & \mathbf{i}^{-(2 n-2)} \omega^{(2 n-2) k} \\
\vdots & \vdots & \vdots & \vdots \\
\mathbf{i}^{(2 n-1)} \omega^{-(2 n-1) k} & \mathbf{i}^{(2 n-2)} \omega^{-(2 n-2) k} & \cdots & 1
\end{array}\right) .
\end{aligned}
$$

and

Therefore, the transition matrix $H(t)$ of the normal Cayley graph $\operatorname{Cay}\left(V_{8 n}, S\right)$ is given by

$$
\begin{aligned}
H(t) & =\exp \left(-\mathbf{i} t \alpha_{1}\right) E_{1}+\exp \left(-\mathbf{i} t \alpha_{2}\right) E_{2}+\exp \left(-\mathbf{i} t \alpha_{3}\right) E_{3}+\exp \left(-\mathbf{i} t \alpha_{4}\right) E_{4} \\
& +\exp \left(-\mathbf{i} t \alpha_{5}\right) E_{5}+\exp \left(-\mathbf{i} t \alpha_{6}\right) E_{6}+\exp \left(-\mathbf{i} t \alpha_{7}\right) E_{7}+\exp \left(-\mathbf{i} t \alpha_{8}\right) E_{8} \\
& +\sum_{j=1}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(E_{j}^{(1)}+E_{j}^{(2)}+E_{j}^{(3)}+E_{j}^{(4)}\right)+\sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(F_{k}^{(1)}+F_{k}^{(2)}+F_{k}^{(3)}+F_{k}^{(4)}\right) .
\end{aligned}
$$

Now we compute the $(u, v)$-th element of the transition matrix. The following three cases arise.
Case 1. If $u \in V_{1}, v \in V_{2}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)+\exp \left(-\mathbf{i} t \alpha_{2}\right) \mathbf{i}^{u-v+1}+(-1)^{u+v+1} \exp \left(-\mathbf{i} t \alpha_{3}\right)+(-\mathbf{i})^{u-v+1} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right. \\
& \left.-\exp \left(-\mathbf{i} t \alpha_{5}\right)+\mathbf{i}^{u-v+3} \exp \left(-\mathbf{i} t \alpha_{6}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{7}\right)+(-\mathbf{i})^{u-v+3} \exp \left(-\mathbf{i} t \alpha_{8}\right)\right)
\end{aligned}
$$

This implies that $\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 8=\frac{1}{n}<1$.
Simillarly, it can be shown that for $u \in V_{1}, v \in V_{4}$ or $u \in V_{2}, v \in V_{1}$ or $u \in V_{2}, v \in V_{3}$ or $u \in V_{3}, v \in V_{2}$ or $u \in V_{3}, v \in V_{4}$ or $u \in V_{4}, v \in V_{1}$ or $u \in V_{4}, v \in V_{3},\left|H(t)_{u v}\right|<1$.
Therefore, in this case PST cannot occur between the vertices $u$ and $v$.
Case 2. If $u, v \in V_{1}$ or $u, v \in V_{2}$ or $u, v \in V_{3}$ or $u, v \in V_{4}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)+\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right)+(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right. \\
& \left.+\exp \left(-\mathbf{i} t \alpha_{5}\right)+\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{6}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{7}\right)+(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{8}\right)\right) \\
& +\frac{1}{4 n} \sum_{j=1}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(\omega^{(u-v) j}+\omega^{-(u-v) j}\right)+\frac{1}{4 n} \sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(\mathbf{i}^{u-v} \omega^{(u-v) k}+\mathbf{i}^{u-v} \omega^{-(u-v) k}\right)
\end{aligned}
$$

This implies that $\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 8+\frac{1}{4 n} \times(2 n-2)+\frac{1}{4 n} \times(2 n-2)=1$.
Therefore, $\left|H(t)_{u v}\right| \leq 1$. Thus $\left|H(t)_{u v}\right|=1$ if and only if for $1 \leq j \leq n-1$ and $1 \leq k \leq n-1$, it holds
that,

$$
\begin{aligned}
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{2}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{4}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\exp \left(-\mathbf{i} t \alpha_{5}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{6}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{7}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{8}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{-(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\mathbf{i}^{u-v} \omega^{(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\mathbf{i}^{u-v} \omega^{-(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right)
\end{aligned}
$$

From the last two equations, since $\omega$ is the $2 n$-th root of unity, we get that $n$ divides $u-v$. Without loss of generality, we assume that $u>v$. This implies that $u-v=n$. Let $t=2 \pi T$. Then the preceding equations implies that

$$
\begin{align*}
& \left(\alpha_{1}-\alpha_{2}\right) T-\frac{n}{4} \in \mathbb{Z}  \tag{22}\\
& \left(\alpha_{1}-\alpha_{3}\right) T \in \mathbb{Z}  \tag{23}\\
& \left(\alpha_{1}-\alpha_{4}\right) T-\frac{n}{4} \in \mathbb{Z}  \tag{24}\\
& \left(\alpha_{1}-\alpha_{5}\right) T \in \mathbb{Z}  \tag{25}\\
& \left(\alpha_{1}-\alpha_{6}\right) T-\frac{n}{4} \in \mathbb{Z}  \tag{26}\\
& \left(\alpha_{1}-\alpha_{7}\right) T \in \mathbb{Z}  \tag{27}\\
& \left(\alpha_{1}-\alpha_{8}\right) T-\frac{n}{4} \in \mathbb{Z}  \tag{28}\\
& \left(\alpha_{1}-\beta_{j}\right) T-\frac{j}{2} \in \mathbb{Z}  \tag{29}\\
& \left(\alpha_{1}-\gamma_{k}\right) T-\frac{n}{4}-\frac{k}{2} \in \mathbb{Z} \tag{30}
\end{align*}
$$

Since $0=\operatorname{tr}(A)=\alpha_{1}+\alpha_{2}+\alpha_{3}+\alpha_{4}+\alpha_{5}+\alpha_{6}+\alpha_{7}+\alpha_{8}+4 \sum_{j=1}^{n-1} \beta_{j}+4 \sum_{k=1}^{n-1} \gamma_{k}$, we have that $8 n \alpha_{1} T \in \mathbb{Z}$, and since $\alpha_{1}=|S|$ is a positive integer, we have $T \in \mathbb{Q}$. This implies that all the eigenvalues of the graph are rational numbers. It is well known that any rational eigenvalue of a graph is an integer. Therefore, in this case the graph is integral.
Our next discussion is distinguished into two parts according as $n \equiv 0(\bmod 4)$ or $n \equiv 2(\bmod 4)$.
(i) $n \equiv 0(\bmod 4)$. In this case, conditions (22) to (30) turns into

$$
\begin{aligned}
& \left(\alpha_{1}-\alpha_{2}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{3}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{4}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{5}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{6}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{7}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{8}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\beta_{j}\right) T-\frac{j}{2} \in \mathbb{Z} \\
& \left(\alpha_{1}-\gamma_{k}\right) T-\frac{k}{2} \in \mathbb{Z}
\end{aligned}
$$

The preceding conditions implies that, $\Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}-1}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}-1}\right)$ are the same, say $\mu_{1}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right), \Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}}\right)$ are all strictly greater than $\mu_{1}$, for $1 \leq j^{\prime} \leq \frac{n-1}{2}$ and $1 \leq k^{\prime} \leq \frac{n-1}{2}$.
(ii) $n \equiv 2(\bmod 4)$. In this case, conditions (22) to (30) turns into

$$
\begin{aligned}
& \left(\alpha_{1}-\alpha_{2}\right) T-\frac{1}{2} \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{3}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{4}\right) T-\frac{1}{2} \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{5}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{6}\right) T-\frac{1}{2} \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{7}\right) T \in \mathbb{Z} \\
& \left(\alpha_{1}-\alpha_{8}\right) T-\frac{1}{1} \in \mathbb{Z} \\
& \left(\alpha_{1}-\beta_{j}\right) T-\frac{j}{2} \in \mathbb{Z} \\
& \left(\alpha_{1}-\gamma_{k}\right) T-\frac{1}{2}-\frac{k}{2} \in \mathbb{Z}
\end{aligned}
$$

The preceding conditions implies that, $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right), \Upsilon_{2}\left(\alpha_{1}-\right.$ $\left.\beta_{2 j^{\prime}-1}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}}\right)$ are the same, say $\mu_{2}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right), \Upsilon_{2}\left(\alpha_{1}-\beta_{2 j^{\prime}}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{2 k^{\prime}-1}\right)$ are all strictly greater than $\mu_{2}$, for $1 \leq j^{\prime} \leq \frac{n-1}{2}$ and $1 \leq k^{\prime} \leq \frac{n-1}{2}$.

Case 3. If $u \in V_{1}, v \in V_{3}$ or $u \in V_{2}, v \in V_{4}$ or $u \in V_{3}, v \in V_{1}$ or $u \in V_{4}, v \in V_{2}$, then

$$
\begin{aligned}
H(t)_{u v} & =\frac{1}{8 n}\left(\exp \left(-\mathbf{i} t \alpha_{1}\right)-\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{2}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right)-(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{4}\right)\right. \\
& \left.+\exp \left(-\mathbf{i} t \alpha_{5}\right)-\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{6}\right)+(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{7}\right)-(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{8}\right)\right) \\
& +\frac{1}{4 n} \sum_{j=1}^{n-1} \exp \left(-\mathbf{i} t \beta_{j}\right)\left(\omega^{(u-v) j}+\omega^{-(u-v) j}\right)+\frac{1}{4 n} \sum_{k=1}^{n-1} \exp \left(-\mathbf{i} t \gamma_{k}\right)\left(-\mathbf{i}^{u-v} \omega^{(u-v) k}-\mathbf{i}^{u-v} \omega^{-(u-v) k}\right)
\end{aligned}
$$

This implies that $\left|H(t)_{u v}\right| \leq \frac{1}{8 n} \times 8+\frac{1}{4 n} \times(2 n-2)+\frac{1}{4 n} \times(2 n-2)=1$.
Therefore, $\left|H(t)_{u v}\right| \leq 1$. Thus $\left|H(t)_{u v}\right|=1$ if and only if for $1 \leq j \leq n-1$ and $1 \leq k \leq n-1$, it holds that,

$$
\begin{aligned}
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{2}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{3}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{4}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\exp \left(-\mathbf{i} t \alpha_{5}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-\mathbf{i}^{u-v} \exp \left(-\mathbf{i} t \alpha_{6}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=(-1)^{u+v} \exp \left(-\mathbf{i} t \alpha_{7}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-(-\mathbf{i})^{u-v} \exp \left(-\mathbf{i} t \alpha_{8}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=\omega^{-(u-v) j} \exp \left(-\mathbf{i} t \beta_{j}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-\mathbf{i}^{u-v} \omega^{(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right) \\
& \exp \left(-\mathbf{i} t \alpha_{1}\right)=-\mathbf{i}^{u-v} \omega^{-(u-v) k} \exp \left(-\mathbf{i} t \gamma_{k}\right)
\end{aligned}
$$

From the last two equations, we get that $u-v=4 n$. Let $t=2 \pi T$. Then the preceding equations implies that

$$
\begin{align*}
& \left(\alpha_{1}-\alpha_{2}\right) T-\frac{1}{2} \in \mathbb{Z}  \tag{31}\\
& \left(\alpha_{1}-\alpha_{3}\right) T \in \mathbb{Z}  \tag{32}\\
& \left(\alpha_{1}-\alpha_{4}\right) T-\frac{1}{2} \in \mathbb{Z}  \tag{33}\\
& \left(\alpha_{1}-\alpha_{5}\right) T \in \mathbb{Z}  \tag{34}\\
& \left(\alpha_{1}-\alpha_{6}\right) T-\frac{1}{2} \in \mathbb{Z}  \tag{35}\\
& \left(\alpha_{1}-\alpha_{7}\right) T \in \mathbb{Z}  \tag{36}\\
& \left(\alpha_{1}-\alpha_{8}\right) T-\frac{1}{2} \in \mathbb{Z} \tag{37}
\end{align*}
$$

$$
\begin{align*}
& \left(\alpha_{1}-\beta_{j}\right) T \in \mathbb{Z}  \tag{38}\\
& \left(\alpha_{1}-\gamma_{k}\right) T-\frac{1}{2} \in \mathbb{Z} \tag{39}
\end{align*}
$$

Since $0=\operatorname{tr}(A)=\alpha_{1}+\alpha_{2}+\alpha_{3}+\alpha_{4}+\alpha_{5}+\alpha_{6}+\alpha_{7}+\alpha_{8}+4 \sum_{j=1}^{n-1} \beta_{j}+4 \sum_{k=1}^{n-1} \gamma_{k}$, we have that $8 n \alpha_{1} T \in \mathbb{Z}$, and since $\alpha_{1}=|S|$ is a positive integer, we have $T \in \mathbb{Q}$. This implies that all the eigenvalues of the graph are rational numbers. It is well known that any rational eigenvalue of a graph is an integer. Therefore, in this case the graph is integral.

The preceding conditions implies that, $\Upsilon_{2}\left(\alpha_{1}-\alpha_{2}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{4}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{6}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{8}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\gamma_{k}\right)$ are the same, say $\mu_{3}$, and $\Upsilon_{2}\left(\alpha_{1}-\alpha_{3}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{5}\right), \Upsilon_{2}\left(\alpha_{1}-\alpha_{7}\right)$ and $\Upsilon_{2}\left(\alpha_{1}-\beta_{j}\right)$ are all strictly greater than $\mu_{3}$, for $1 \leq j \leq n-1$ and $1 \leq k \leq n-1$.
Using the same technique in Section 3.1, we can show that the minimum time at which $\Gamma$ has PST between $u$ and $v$ is $\frac{\pi}{M}$, where $M=\operatorname{gcd}\left(\alpha-\alpha_{1}: \alpha \in \operatorname{Spec}(\Gamma) \backslash\left\{\alpha_{1}\right\}\right)$.
This completes the proof.

## References

[1] R.J. Angeles-Canul, R. Norton, M. Opperman, C. Paribello, M. Russell and C. Tamon. Perfect state transfer, integral circulants and join of graphs. Quantum. Inf. Comput. 10 (2010), 325-342.
[2] M. Arezoomand, F. Shafiei and M. Ghorbani. Perfect state transfer on Cayley graphs over the dicyclic group. Linear Algebra Appl. 639 (2022), 116-134.
[3] M. Basic and M.D. Petkovic. Perfect state transfer in integral circulant graphs of non-square-free order. Linear Algebra Appl. 433(1) (2010), 149-163.
[4] M. Basic, M.D. Petkovic and D. Stevanovic. Perfect state transfer in integral circulant graphs. Appl. Math. Lett. 22(7) (2009), 1117-1121.
[5] M. Basic and M.D. Petkovic. Some classes of integral circulant graphs either allowing or not allowing perfect state transfer. Appl. Math. Lett. 22(10) (2009), 1609-1615.
[6] A. Bernasconi, C. Godsil and S. Severini. Quantum networks on cubelike graphs. Phys. Rev. A. (3) $78(5)(2008), 052320$.
[7] S. Bose. Quantum communication through an unmodulated spin chain. Phys. Rev. lett. 91(20) (2003), 207901.
[8] X. Cao and K. Feng. Perfect state transfer on Cayley graphs over dihedral groups. Linear Multilinear Algebra. 69(2) (2021), 343-360.
[9] W. Cheung and C. Godsil. Perfect state transfer in cubelike graphs. Linear Algebra Appl. 435(10) (2011), 2468-2474.
[10] M. Christandl, N. Datta, T. Dorlas, A. Ekert, A. Kay and A. Landahl. Perfect transfer of arbitrary states in quantum spin networks. Phys. Rev. A. 71 (2005), 032312.
[11] G. Coutinho and C. Godsil. Perfect state transfer in products and covers of graphs. Linear Multilinear Algebra. 64(2) (2016), 235-246.
[12] G. Coutinho. Quantum state transfer in graphs [PhD dissertation]. University of Waterloo, (2014).
[13] M. R. Darafsheh and N. S. Poursalavati. On the existence of the orthogonal basis of the symmetry classes of tensors associated with certain groups. SUT J. Math. 37(1) (2001), 1-17.
[14] E. Farhi and S. Gutmann. Quantum computation and decision trees. Phys. Rev. A. (3) 58 (1998), 915-928.
[15] C. Godsil and G. Royle. Algebraic Graph Theory. Graduate Texts in Mathematics. Vol. 207. Springer-Verlag, New York, (2001).
[16] C. Godsil. Periodic graphs. Electron. J. Comb. 18(1) (2011), \#P23.
[17] C. Godsil. State transfer on graphs. Discrete Math. 312(1) (2012), 129-147.
[18] C. Godsil. When can perfect state transfer occur? Electron. J. Linear Algebra. 23 (2012), 877-890.
[19] G. James and M. Liebeck. Representations and Characters of Groups. Second Edition. Cambridge University Press, New York, (2001).
[20] A. Khalilipour and M. Ghorbani. Perfect state transfer on Cayley graphs over groups. J. Disc. Math. Appl. 8(4) (2023), 213-226.
[21] G. Luo, X. Cao, D. Wang and X. Wu. Perfect quantum state transfer on Cayley graphs over semidihedral groups. Linear Multilinear Algebra. 70(21) (2022), 6358-6374.
[22] H. Pal and B. Bhattacharjya. Perfect state transfer on NEPS of the path on three vertices. Discrete Math. 339 (2016), 831-838.
[23] M.D. Petkovic and M. Basic. Further results on the perfect state transfer in integral circulant graphs. Comput. Math. Appl. 61 (2011), 300-312.
[24] B. Steinberg. Representation Theory of Finite Groups. Universitext, Springer, New York, 2012.
[25] T. Tan, K. Feng and X. Cao. Perfect state transfer on abelian Cayley graphs. Linear Algebra Appl. 563 (2019), 331-352.

